

Mass Formula (Even-odd mass difference)

Extra binding when like nucleons form a spin-zero pair

Example:

	Binding energy (MeV)
${}^{210}_{82}\text{Pb}_{128} = {}^{208}_{82}\text{Pb}_{126} + 2n$	1646.6
${}^{210}_{83}\text{Bi}_{127} = {}^{208}_{82}\text{Pb}_{126} + n + p$	1644.8
${}^{209}_{82}\text{Pb}_{127} = {}^{208}_{82}\text{Pb}_{126} + n$	1640.4
${}^{209}_{83}\text{Bi}_{126} = {}^{208}_{82}\text{Pb}_{126} + p$	1640.2

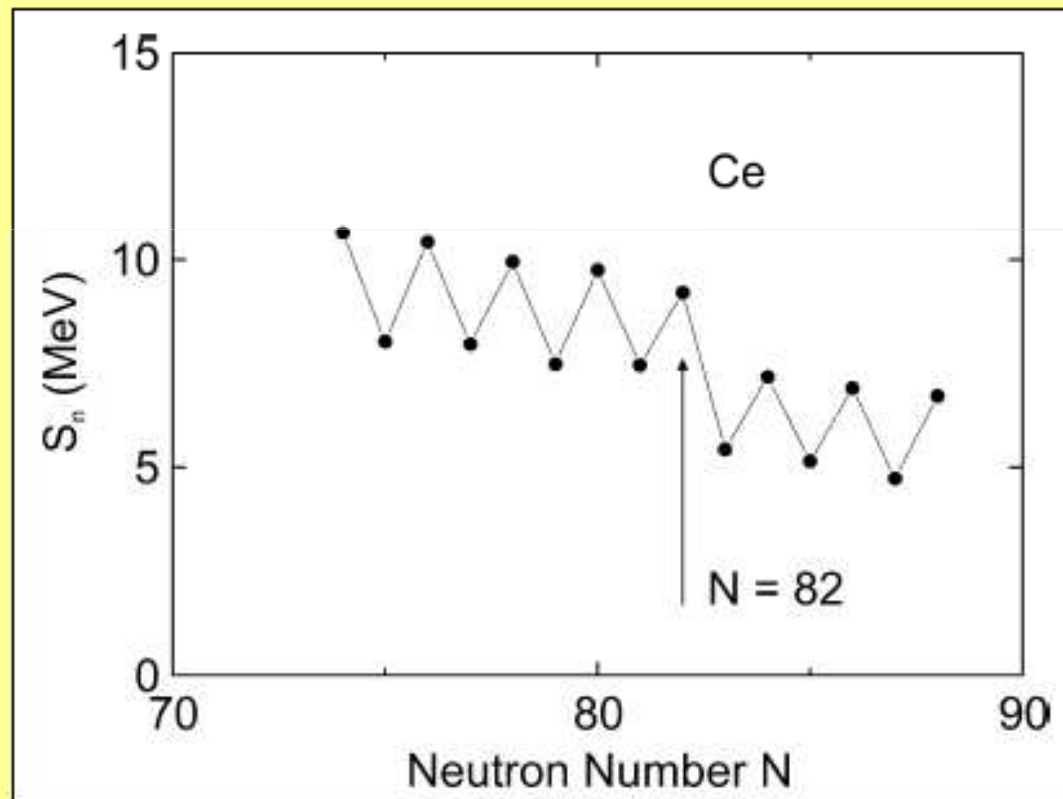
$$\begin{aligned} B_{\text{pair}} &= \Delta && \text{(for even - even)} \\ &= 0 && \text{(for even - odd)} \\ &= -\Delta && \text{(for odd - odd)} \end{aligned}$$

More later

Evidence for Pairing Correlations in Nuclei

- (i) Odd-even effect: mass of an odd-even nucleus is larger than the mean of adjacent two even-even nuclear masses \rightarrow shows up in S_n and S_p for all nuclei.

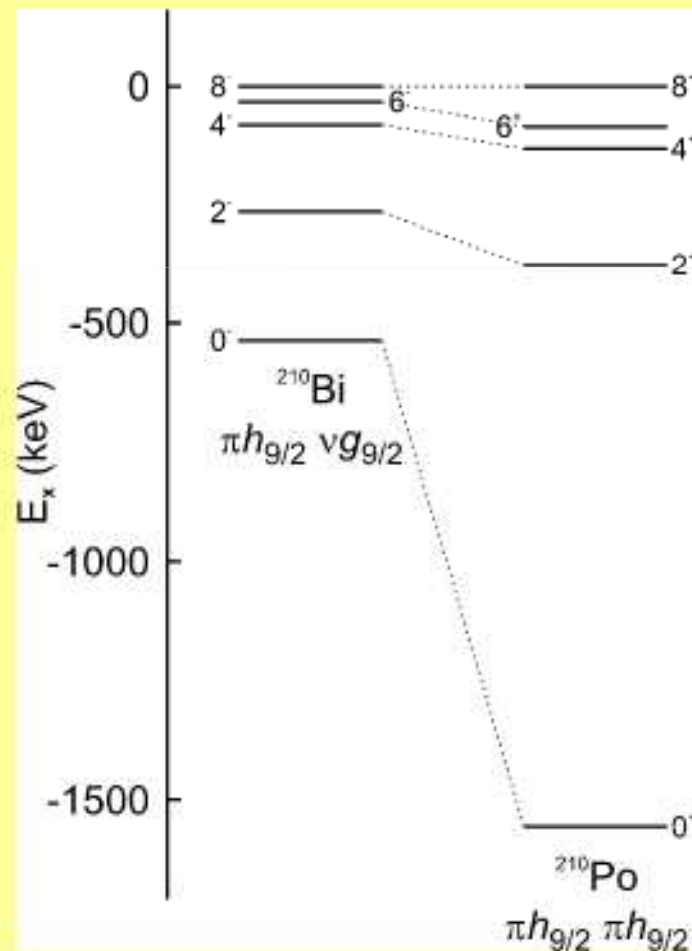
Example: $S_n = BE(A, Z) - BE(A-1, Z)$ of Ce nuclei



- Behavior points towards pair formation of nucleons.

Evidence for pairing correlations in nuclei

- (ii) Energy spectra for nuclei near closed shells ($A \pm 2$, $A \pm 4$) show a clear gap for the 0^+ g.s.



- In ^{210}Po the configuration outside the doubly-closed shell core of ^{208}Pb is $(1h_{9/2})^2$. If there were no interaction between the two π 's constituting the pair, i.e. if they behaved like independent particles, the various $(1h_{9/2})^2$ spin couplings, which reflect the orbital alignments, would lead to states degenerate in energy.

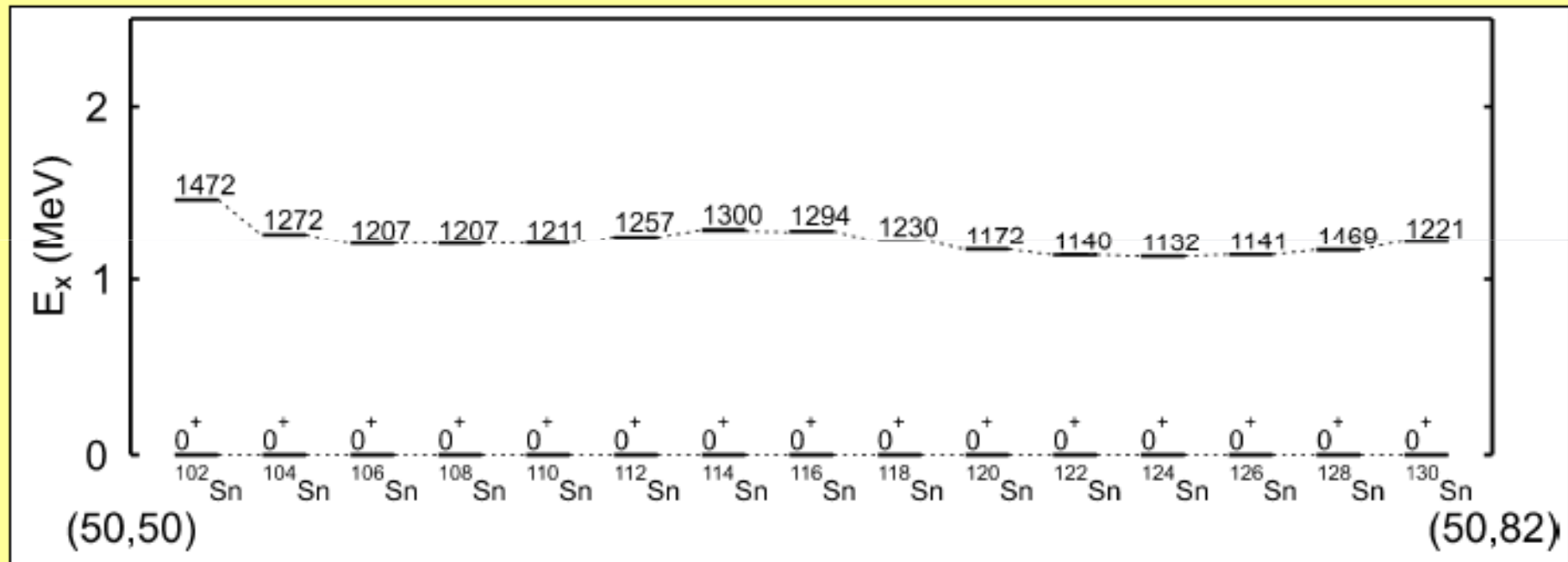
→ correlated pair of two π 's

- Pairing effect $\approx 2\Delta$

Evidence for Pairing Correlations in Nuclei

(ii) The excitation energy of the first excited 2^+ state in nuclei remains remarkably constant over large intervals of neutron (proton) numbers.

Example: 2^+_1 excitation energy in Sn nuclei

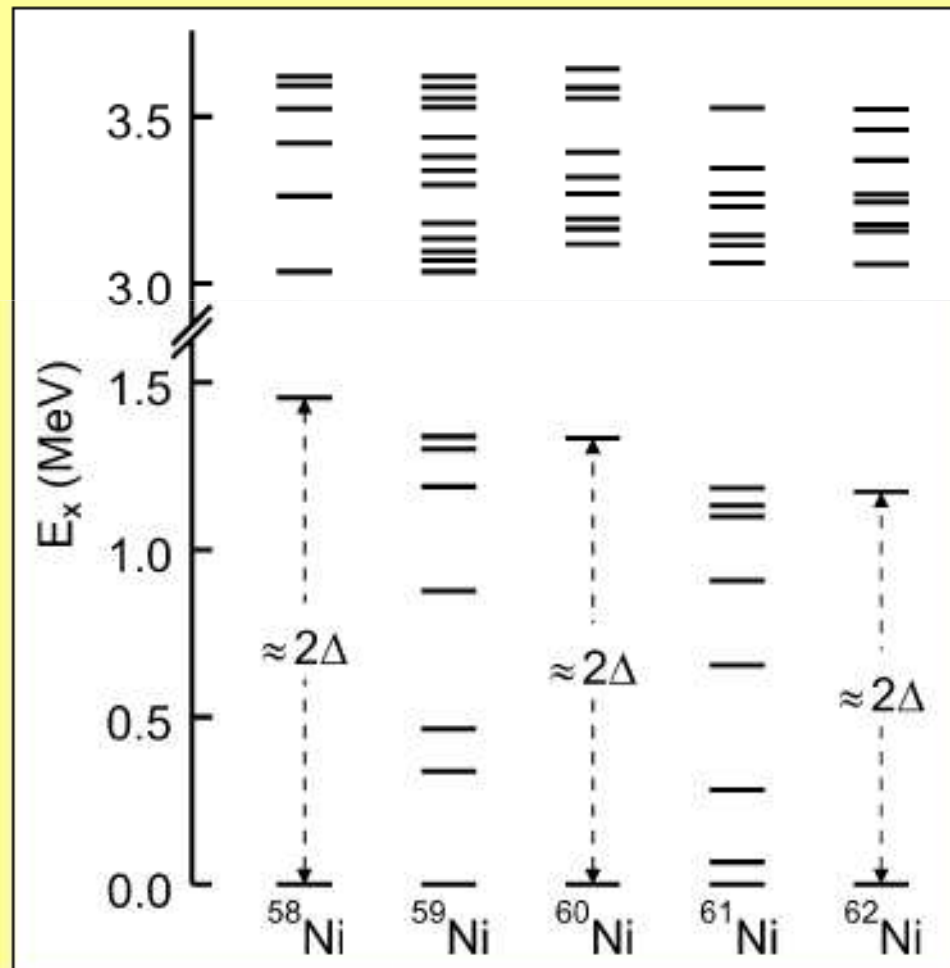


- These 2^+ states are not rotational states but are connected to a coherent pairing condensate.
- Pair breaking energy: $2\Delta \approx 2$ MeV

Evidence for Pairing Correlations in Nuclei

(iii) Energy gap: odd-even and even-odd nuclei (especially deformed nuclei) have energy spectra different from even-even nuclei.

Example: Ni isotopes



- e-e nuclei: only a few states at most (vibrations, rotations) appear below the pairing gap 2Δ .

- But in o-e and e-o nuclei (where the last nucleon is unpaired) many s.p. and collective states appear.

- Note: above the pair breaking energy 2Δ many excited states are possible \rightarrow level density $\rho = \rho(\Delta)$

\rightarrow **2nd Lecture**

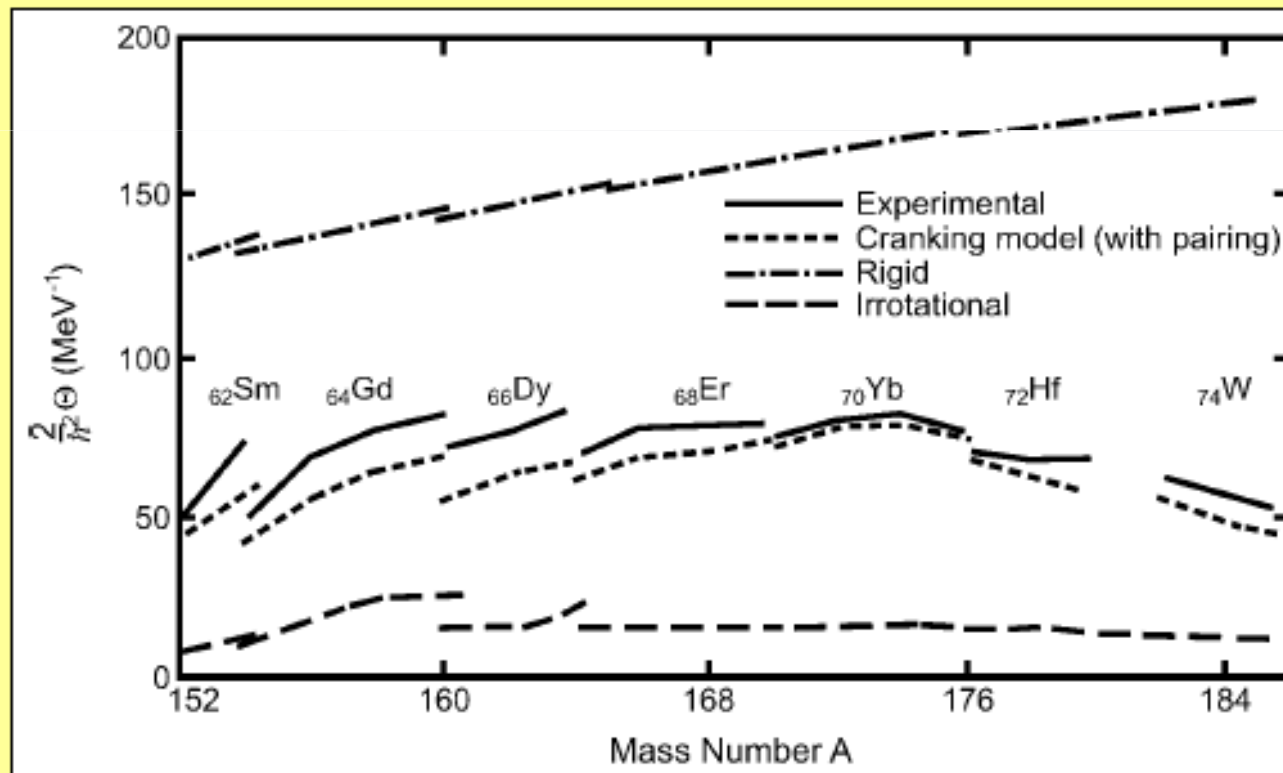
Evidence for Pairing Correlations in Nuclei

(iv) Moment of inertia: extracted from level spacing in rotational bands

$$E = \frac{\hbar^2}{2\Theta} J(J + 1)$$

deviates about a factor of two from the rigid rotor values.

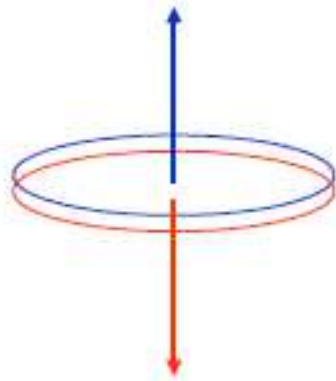
Example: Moments of inertia of even-nuclei in the rare-earth region



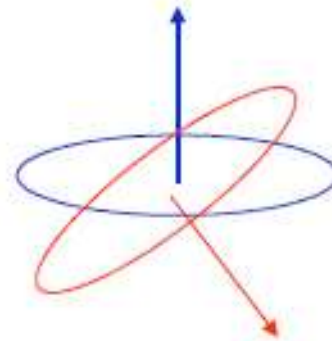
● $\Theta_{irrot} < \Theta_{exp} < \Theta_{rigid}$

● Pairing correlations have a dramatic influence on collective modes.

Simple interpretation:



$I=0$ pair



$I \neq 0$ pair

The spatial overlap is the largest for the $I=0$ pair.

“Pairing Correlation”

(note) The $I=2j$ pair is unfavoured due to the Pauli principle.

(note)

$$\psi(l^2; L=0) = \sum_{\mu} \langle l\mu l-\mu | L=0, 0 \rangle Y_{l\mu}(\hat{r}_1) Y_{l-\mu}(\hat{r}_2) = Y_{l0}(\theta_{12}) / \sqrt{4\pi}$$