

# Nuclear beta decay

- Introduction
- Energy release
- Fermi theory of  $\beta$ -decay
- Shape of  $\beta$  spectrum (Kurie plot)
- Total decay rate
- Selection rules
- Neutrino mass
- Summary

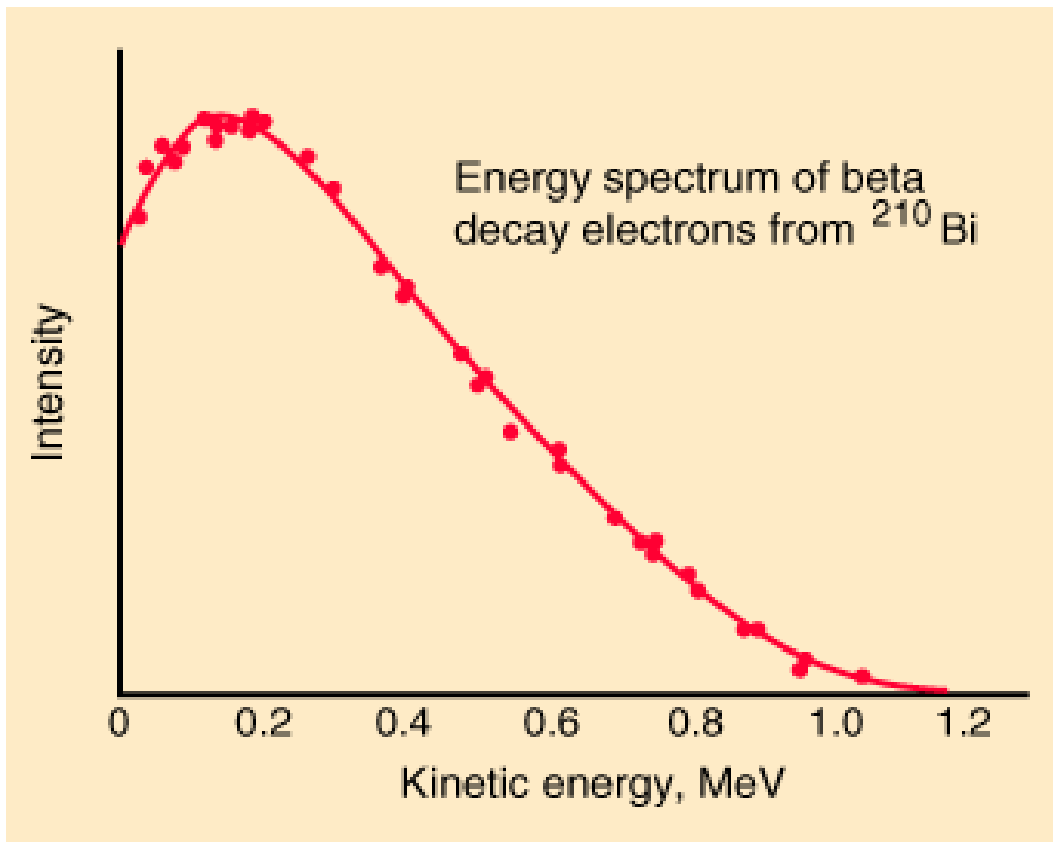
# Introduction

Basic process: conversion of a proton to a neutron or of a neutron to a proton.

Both  $Z$  and  $N$  change by one unit:  $Z \rightarrow Z \pm 1$ ,  $N \rightarrow N \mp 1$ , so that  $A = Z + N$  remains constant  $\Rightarrow$  slide down the mass parabola.

- $\beta^-$ : emission of an  $e^-$
- $\beta^+$ : emission of a  $e^+$
- **EC**: capture of an orbital electron (in competition with  $\beta^+$ )

In case of  $\beta$  emission, continuous energy spectrum:



# Energy release

- $\beta^-$

$$M^*(Z, A) + Zm_e = (M^*(Z+1, A) + (Z+1)m_e) + E_0$$

$$M(Z, A) = M(Z+1, A) + E_0$$

$$Q_{\beta^-} = M(Z, A) - M(Z+1, A) = E_0$$

- $\beta^+$

$$M^*(Z, A) = (M^*(Z-1, A) + m_e) + E_0$$

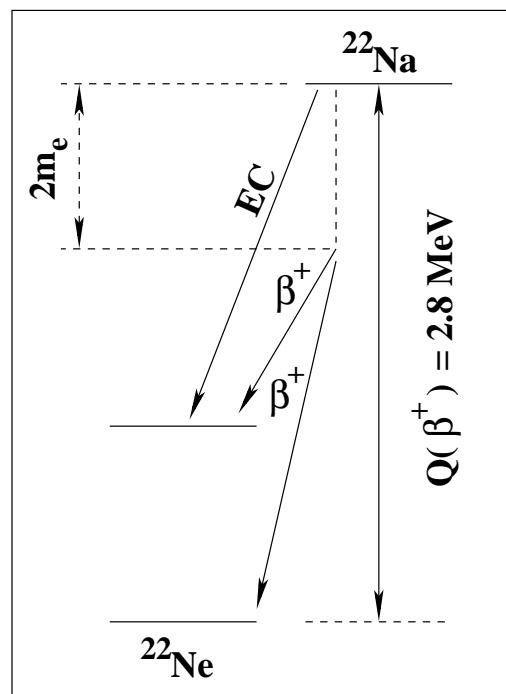
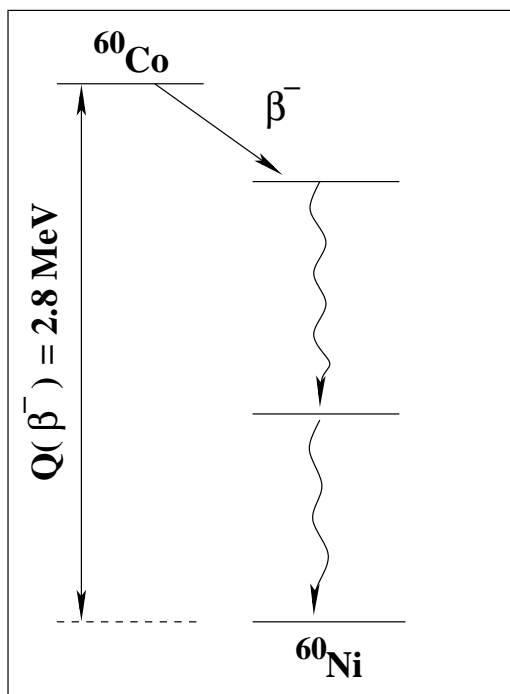
$$M^*(Z, A) + Zm_e = M^*(Z-1, A) + (Z-1)m_e + 2m_e + E_0$$

$$M(Z, A) = M(Z-1, A) + 2m_e + E_0$$

$$Q_{\beta^+} = M(Z, A) - M(Z-1, A) = 2m_e + E_0$$

- Electron capture

$$Q_{EC} = M(Z, A) - M(Z+1, A) = E_0 + B_i$$



# Fermi theory of $\beta$ -decay

Treat the decay-causing interaction as a weak perturbation, find a relationship for the transition rate: **Fermi's Golden Rule**

$$\lambda = \frac{2\pi}{\hbar} |V_{fi}|^2 \rho(E_f)$$

- $\lambda$  : total decay rate
- $V_{fi} = g \int [\psi_f^* \phi_e^* \phi_\nu^*] \hat{O} \psi_i dV = g M_{fi}$
- $\rho(E_f) = \frac{dn}{dE_f}$  : density of final states

Taking the  $e^-$  and  $\nu$  wave functions as plane waves, expanding the exponentials and keeping the first term gives the **allowed approximation** ( $e^{i\vec{p}\cdot\vec{r}} \approx 1 + i\vec{p}\cdot\vec{r} + \dots$ ).

Taking into account the nuclear Coulomb field and higher order terms in the exponential expansion yields the partial decay rate:

$$d\lambda = N(p) \propto \underbrace{p^2 (Q - T_e)^2}_{\text{statistical factor}} \underbrace{F(Z', p)}_{\text{Fermi function}} |M_{fi}|^2 \underbrace{S(p, q)}_{\text{forb. corr.}}$$

Including phase space yields (**density of final states**):

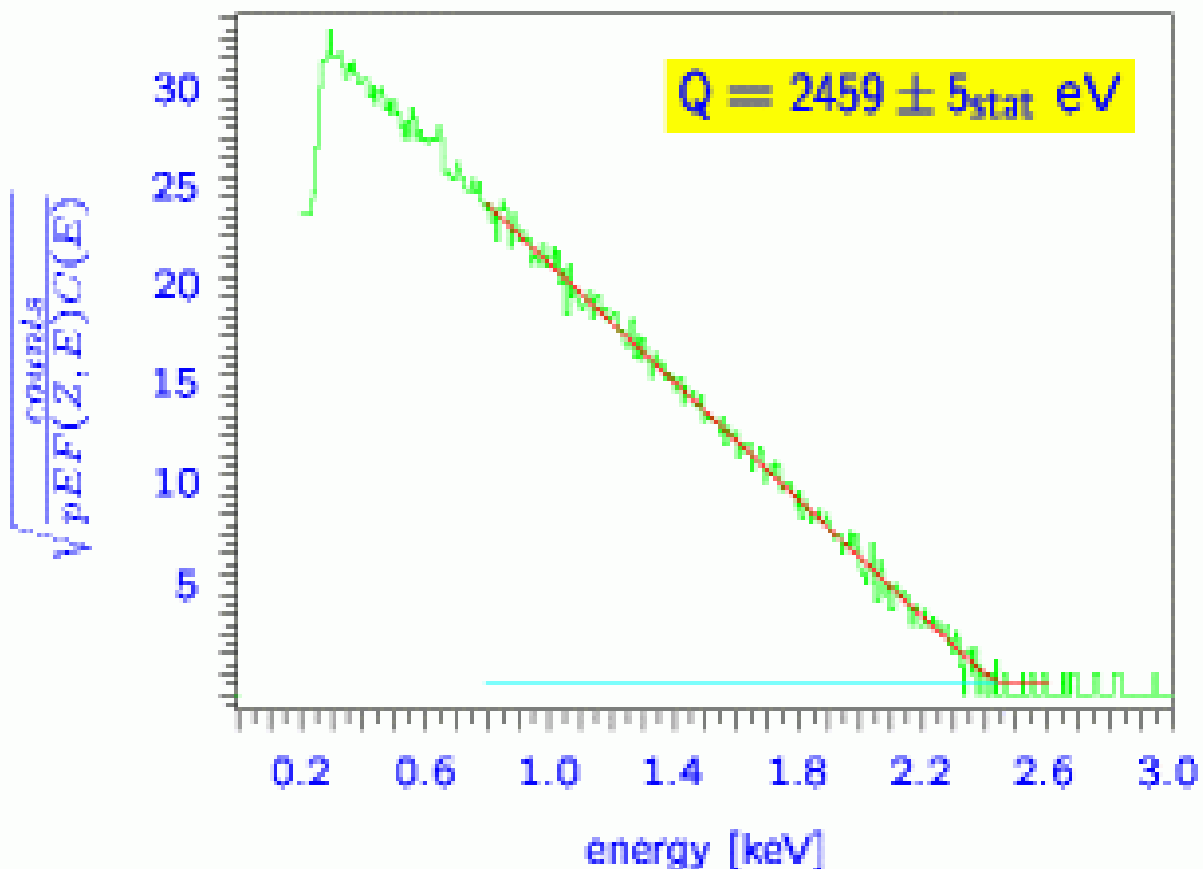
$$d\lambda = \frac{g^2 |M_{fi}|^2}{2\pi \hbar^7 c^3} F(Z', p) p^2 (Q - T_e)^2 dp$$

# Shape of $\beta$ spectrum (Kurie plot)

In the **allowed approximation**, we can rewrite:

$$Q - T_e \propto \sqrt{\frac{N(p)}{p^2 F(Z', p)}}$$

Plotting  $\sqrt{\frac{N(p)}{p^2 F(Z', p)}}$  against  $T_e$  gives a straight line intercepting the x-axis at the decay energy  $Q \Rightarrow$  **Kurie** (or **Fermi**)-plot.



In the case of forbidden decays, the Kurie plot can be linearized by the correction factor  $S(p, q)$ .

# Total decay rate

To find the total decay rate, integrate over all values of the electron momentum  $p$ .

$$\begin{aligned}\lambda &= \frac{g^2 |M_{fi}|^2}{2\pi\hbar^7 c^3} \int_0^{p_{max}} F(Z', p) p^2 (Q - T_e)^2 dp \\ &= \frac{g^2 |M_{fi}|^2}{2\pi\hbar^7 c^3} f(Z', E_0) \quad \text{and}\end{aligned}$$

$$f(Z', E_0) \propto \int_0^{p_{max}} F(Z', p) p^2 (E_0 - E_e)^2 dp$$

where  $f(Z', E_0)$  only depends on  $Z'$  and the maximum electron total energy  $E_0$  and is called the **Fermi integral**. It is tabulated in *Feenberg and Trigg* diagrams.

We can also write it as:

$$ft_{1/2} = 0.693 \frac{2 \pi^3 \hbar^7}{g^2 m_e^5 c^4 |M_{fi}|^2}$$

$ft_{1/2}$  is called the **Comparative half-life** or **ft value**. This quantity gives an indication on the matrix element  $|M_{fi}|^2$  responsible for the transition. Because the range in  $ft_{1/2}$  is very large, what is usually quoted is the value of  $\log_{10}(ft)$  with  $t$  given in seconds.

Depending on the **selection rules** one obtains typical ft-values.

# Selection rules

- **Allowed and superallowed transitions:**

$\Delta\pi = \text{NO}$ ,  $\Delta J = 0$  **Fermi** (singulet)

$\Delta\pi = \text{NO}$ ,  $\Delta J = 0, \pm 1$ , (no  $0 \rightarrow 0$ ) **Gamov-Teller** (triplet)

superallowed:  $\log_{10} ft \approx 3.5$  mirror decays

allowed :  $\log_{10} ft \approx 4 - 7$

Example:  $n \rightarrow p$ ,  $^{14}\text{O} \rightarrow ^{14}\text{N}$

- **Once forbidden** (unique if  $\Delta J = \pm 2$ )

$\Delta\pi = \text{YES}$ ,  $\Delta J = 0, \pm 1, \pm 2$

$\log_{10} ft \approx 6 - 10$

Example:  $^{17}\text{N} \rightarrow ^{17}\text{O} (\frac{1}{2}^- \rightarrow \frac{5}{2}^+)$

- **Twice forbidden** (unique)

$\Delta\pi = \text{NO}$ ,  $\Delta J = \pm 2, \pm 3$

$\log_{10} ft \approx 12 - 14$

Example:  $^{137}\text{Cs} \rightarrow ^{137}\text{Ba} (\frac{7}{2}^+ \rightarrow \frac{3}{2}^+)$

Forbidden decays are less probable because they contain an orbital angular momentum change.

Example:

$$E_0 = 1 \text{ MeV}, p_{max}(e^-) = 1.4 \text{ MeV}/c, R = 3 \text{ fm}$$

$$\Rightarrow pR = 4.2 \text{ fm MeV}/c = 0.2 \hbar < 1 \hbar$$

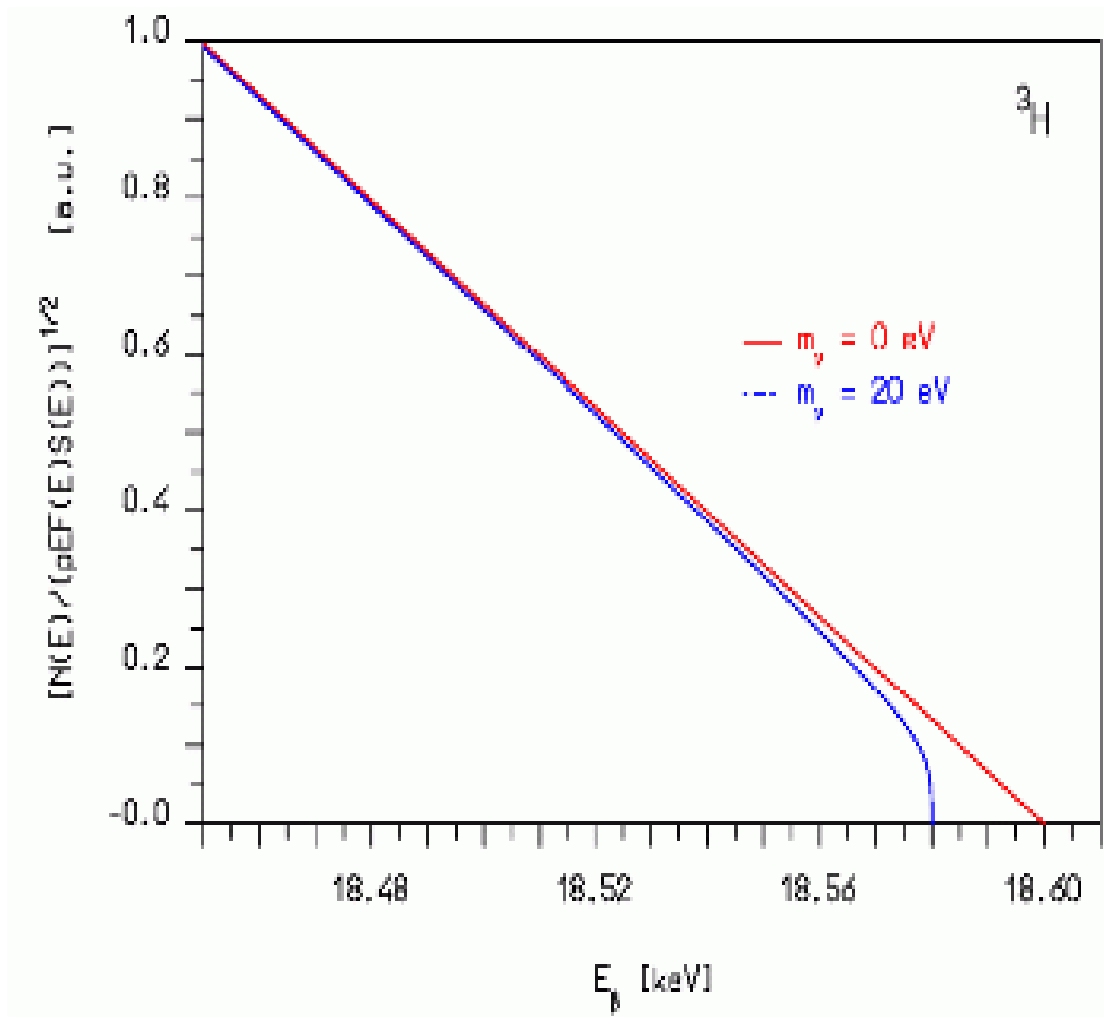
# Neutrino mass

Looking at the endpoint energy of the  $\beta$ -spectrum is a method to determine the neutrino mass (upper limit is now 17 eV).

Taking the neutrino mass into account yields for the decay rate:

$$d\lambda = \frac{g^2 |M_{fi}|^2}{2\pi\hbar^7 c^3} F(Z', p) p^2 (Q - T_e)^2 \sqrt{1 - \frac{m_\nu^2 c^4}{(Q - T_e)^2}} dp$$

At the endpoint ( $Q - T_e = 0$ ), if  $m_\nu = 0$  then  $\frac{d\lambda}{dp} \rightarrow 0$ . If  $m_\nu \neq 0$  then  $\frac{d\lambda}{dp} \rightarrow \infty$ .





# Summary

- Based on Fermi's Golden Rule (perturbation theory)
- $\beta$ -spectrum is a continuous spectrum (3 body process)
- $\log_{10}(ft)$ -values allows one to predict the spin and parity of nuclear states
- Powerfull spectroscopic tool (but complicated)