

Time-dependent Green's Functions approach to nuclear reactions

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Encuentros de Física Nuclear 2010

A. Rios et al., arXiv:1009.0215

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Outline



1 Motivation

- 2 1D mean-field dynamics
- 3 Correlated time evolution





Time matters!



Nuclear reactions are time-dependent processes

- · Advancements of time-dependent many-body techniques are needed
 - Central collisions of heavy isotopes \Rightarrow many participants, rearrangement
 - Low-energy fusion reactions ⇒ sub-barrier fusion, neck formation
 - Response of finite nuclei \Rightarrow collective phenomena, deexcitation
- · Properties of time-dependent Green's function's approach
 - Fully quantal approach
 - Consistent many-body framework
 - · Correlations beyond the mean-field approximation







Kadanoff-Baym equations



$$\begin{cases} i\frac{\partial}{\partial t_1} + \frac{\nabla_1^2}{2m} \end{bmatrix} \mathcal{G}^{\lessgtr}(\mathbf{1}\mathbf{1}') = \int \mathrm{d}\bar{\mathbf{r}}_1 \Sigma_{HF}(\mathbf{1}\bar{\mathbf{1}}) \mathcal{G}^{\lessgtr}(\bar{\mathbf{1}}\mathbf{1}') \\ + \int_{t_0}^{t_1} \mathrm{d}\bar{\mathbf{1}} \left[\Sigma^{>}(\mathbf{1}\bar{\mathbf{1}}) - \Sigma^{<}(\mathbf{1}\bar{\mathbf{1}}) \right] \mathcal{G}^{\lessgtr}(\bar{\mathbf{1}}\mathbf{1}') - \int_{t_0}^{t_1'} \mathrm{d}\bar{\mathbf{1}} \Sigma^{\lessgtr}(\mathbf{1}\bar{\mathbf{1}}) \left[\mathcal{G}^{>}(\bar{\mathbf{1}}\mathbf{1}') - \mathcal{G}^{<}(\bar{\mathbf{1}}\mathbf{1}') \right] \end{cases}$$

$$\begin{split} \left\{ -i\frac{\partial}{\partial t_{1'}} + \frac{\nabla_{1'}^2}{2m} \right\} \mathcal{G}^{\lessgtr}(\mathbf{11}') &= \int \mathrm{d}\bar{\mathbf{r}}_1 \mathcal{G}^{\lessgtr}(\mathbf{1}\bar{\mathbf{1}}) \Sigma_{HF}(\bar{\mathbf{11}}') \\ &+ \int_{t_0}^{t_1} \mathrm{d}\bar{\mathbf{1}} \left[\mathcal{G}^{\gt}(\mathbf{1}\bar{\mathbf{1}}) - \mathcal{G}^{<}(\mathbf{1}\bar{\mathbf{1}}) \right] \Sigma^{\lessgtr}(\bar{\mathbf{11}}') - \int_{t_0}^{t_1'} \mathrm{d}\bar{\mathbf{1}} \, \mathcal{G}^{\lessgtr}(\mathbf{1}\bar{\mathbf{1}}) \left[\Sigma^{\gt}(\bar{\mathbf{11}}') - \Sigma^{<}(\bar{\mathbf{11}}') \right] \end{split}$$

- Evolution of non-equilibrium systems from general principles
- Include correlation and memory effects, via self-energies
- Conservation laws are preserved
- Complicated numerical solution, but very universal framework

Kadanoff & Baym, Quantum Statistical Mechanics (1962)

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Kadanoff-Baym equations $\mathcal{G}^{<}(\mathbf{11}') = i \left\langle \Phi_{0} \middle| \hat{a}^{\dagger}(\mathbf{1}') \hat{a}(\mathbf{1}) \middle| \Phi_{0} \right\rangle \qquad \mathcal{G}^{>}(\mathbf{11}') = -i \left\langle \Phi_{0} \middle| \hat{a}(\mathbf{1}) \hat{a}^{\dagger}(\mathbf{1}') \middle| \Phi_{0} \right\rangle$ $\left\{i\frac{\partial}{\partial t_1} + \frac{\nabla_1^2}{2m}\right\} \mathcal{G}^{\lessgtr}(\mathbf{11}') = \int \mathrm{d}\mathbf{\bar{r}}_1 \Sigma_{HF}(\mathbf{1\bar{1}}) \mathcal{G}^{\lessgtr}(\mathbf{\bar{11}}')$ $+ \int_{t}^{t_1} \! \mathrm{d}\bar{\mathbf{1}} \left[\Sigma^{>}(\mathbf{1}\bar{\mathbf{1}}) - \Sigma^{<}(\mathbf{1}\bar{\mathbf{1}}) \right] \mathcal{G}^{\lessgtr}(\bar{\mathbf{1}}\mathbf{1}') - \int_{t_0}^{t_1'} \! \mathrm{d}\bar{\mathbf{1}} \Sigma^{\lessgtr}(\mathbf{1}\bar{\mathbf{1}}) \left[\mathcal{G}^{>}(\bar{\mathbf{1}}\mathbf{1}') - \mathcal{G}^{<}(\bar{\mathbf{1}}\mathbf{1}') \right]$ $\Sigma_{HF} \Rightarrow$

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$$\Sigma_{HF} \Rightarrow \qquad \langle \hat{O} \rangle = -i \lim_{x \to x'} \int \mathrm{d}x \, o(x) \, \mathcal{G}^{<}(x, x'; 0)$$

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Collisions of 1D slabs





- Frozen & extended y, z coordinates, dynamics in x
- Simple zero-range mean field (1D-3D connection)

$$U(x) = \frac{3}{4}t_0 n(x) + \frac{2+\sigma}{16}t_3 [n(x)]^{(\sigma+1)}$$

- Attempt to understand nuclear Green's functions
- 1D provide a simple visualization
- · Learning before correlations & higher D's

A. Rios et al., arXiv:1009.0215

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Collisions of 1D slabs: fusion



 $E_{CM}/A=0.1\,{\rm MeV}$



Collisions of 1D slabs: fusion $E_{CM}/A = 0.1 \, \text{MeV}$







• Systematic analysis of energy dependence

- Study of origin of off-diagonal elements
- Assessment of elimination schemes
 - Super-operator erasure of off-diagonal elements
 - Coarse-graining of off-diagonal meshes
- Impact on time reversibility

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Wigner distribution







• Fourier transform along relative variable (Wigner transform)

$$f_W(x,p) = \int \frac{\mathrm{d}x_r}{2\pi} \, e^{-ipx_r} \, \mathcal{G}^<\left(x + \frac{x_r}{2}, x - \frac{x_r}{2}\right)$$

- Simultaneous information on real and momentum space
- Lack of correlations \Rightarrow Wigner valley

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Time evolution beyond the mean-field



$$\left\{-i\frac{\partial}{\partial t_{1}}-\frac{\nabla_{1}^{2}}{2m}-\int\!\mathrm{d}\bar{\mathbf{r}}_{1}\Sigma_{HF}(\mathbf{1}\bar{\mathbf{1}})\right\}\mathcal{G}^{\lessgtr}(\mathbf{1}\mathbf{1}')=\underbrace{\int_{t_{0}}^{t_{1}}\!\mathrm{d}\bar{\mathbf{1}}\,\Sigma^{R}(\mathbf{1}\bar{\mathbf{1}})\mathcal{G}^{\lessgtr}(\bar{\mathbf{1}}\mathbf{1}')+\int_{t_{0}}^{t_{1}'}\!\mathrm{d}\bar{\mathbf{1}}\,\Sigma^{\lessgtr}(\mathbf{1}\bar{\mathbf{1}})\mathcal{G}^{A}(\bar{\mathbf{1}}\mathbf{1}')}_{I_{1}^{\lessgtr}(\mathbf{1},\mathbf{1}';t_{0})}$$



• Direct Born approximation ⇒ simplest conserving approximation

FFT to compute convolution integrals

Collision integrals \Rightarrow memory effects in 2D \Rightarrow (t,t')

Time evolution beyond the mean-field



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$$\begin{split} \Sigma^{\lessgtr}(p,t;p',t') &= \int \frac{\mathrm{d}p_1}{2\pi} \frac{\mathrm{d}p_2}{2\pi} V(p-p_1) V(p'-p_2) \mathcal{G}^{\lessgtr}(p_1,t;p_2,t') \Pi^{\lessgtr}(p-p_1,t;p'-p_2,t') \\ \Pi^{\lessgtr}(p,t;p',t') &= \int \frac{\mathrm{d}p_1}{2\pi} \frac{\mathrm{d}p_2}{2\pi} \mathcal{G}^{\lessgtr}(p_1,t;p_2,t') \mathcal{G}^{\gtrless}(p_2-p',t';p_1-p,t) \end{split}$$

- Direct Born approximation \Rightarrow simplest conserving approximation
- FFT to compute convolution integrals
 - Collision integrals \Rightarrow memory effects in 2D \Rightarrow (t, t')

Time evolution beyond the mean-field



$$\left\{-i\frac{\partial}{\partial t_{1}}-\frac{\nabla_{1}^{2}}{2m}-\int \mathrm{d}\bar{\mathbf{r}}_{1}\Sigma_{HF}(\mathbf{1}\bar{\mathbf{1}})\right\}\mathcal{G}^{\lessgtr}(\mathbf{1}\mathbf{1}')=\underbrace{\int_{t_{0}}^{t_{1}}\mathrm{d}\bar{\mathbf{1}}\Sigma^{R}(\mathbf{1}\bar{\mathbf{1}})\mathcal{G}^{\lessgtr}(\bar{\mathbf{1}}\mathbf{1}')+\int_{t_{0}}^{t_{1}'}\mathrm{d}\bar{\mathbf{1}}\Sigma^{\lessgtr}(\mathbf{1}\bar{\mathbf{1}})\mathcal{G}^{A}(\bar{\mathbf{1}}\mathbf{1}')}_{I_{1}^{\lessgtr}(\mathbf{1},\mathbf{1}';t_{0})}$$

$$\begin{split} I_1^{>}(p_1,t_1;p_{1'},t_{1'}) &= \int_{t_0}^{t_1} \!\! \mathrm{d}\bar{t} \int \!\! \frac{\mathrm{d}\bar{p}}{2\pi} \left[\Sigma^{>}(p_1,t_1;\bar{p},\bar{t}) - \Sigma^{<}(p_1,t_1;\bar{p},\bar{t}) \right] \mathcal{G}^{>}(\bar{p},\bar{t};p_{1'},t_{1'}) \\ &- \int_{t_0}^{t_{1'}} \!\! \mathrm{d}\bar{t} \int \!\! \frac{\mathrm{d}\bar{p}}{2\pi} \Sigma^{>}(p_1,t_1;\bar{p},\bar{t}) \left[\mathcal{G}^{<}(\bar{p},\bar{t};p_{1'},t_{1'}) - \mathcal{G}^{>}(\bar{p},\bar{t};p_{1'},t_{1'}) \right] \end{split}$$

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Correlated fermions in a trap









- Mean-field, correlated and off-diagonal elements in 1D \surd
- Lessons learned ⇒ Progressive understanding of higher D
- · Potential for applications in nuclear reactions & structure
- Microscopic understanding of dissipation
- · Response for nuclei including collision width
- Multidisciplinary research: from quantum dots to cosmology!
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Golabek & Simenel, PRL 103, 042701 (2009)

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Piekarewicz & Centelles, PRC 79, 054311 (2009)

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Balzer et al., PRB 79, 245306 (2009)

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Gracias! Gràcies!



A. Rios et al., arXiv:1009.0215



Mean-field evolution: implementation



- · The mean-field is time-local
 - $\Sigma_{HF}(\mathbf{11}') = \delta(t_1 t_{1'}) \Sigma_{HF}(x_1, x_{1'})$
 - Only $t_1 = t_{1'} = t$ elements needed: $\mathcal{G}^{<}(t_1, t_{1'}) \Rightarrow \mathcal{G}^{<}(t)$
- KB equations reduce to one differential equation

$$\begin{split} \frac{\partial}{\partial t} \mathcal{G}^{<}(x, x'; t) &= \left\{ -\frac{1}{2m} \frac{\partial^2}{\partial x^2} + U(x, t) \right\} \mathcal{G}^{<}(x, x'; t) \\ &- \left\{ -\frac{1}{2m} \frac{\partial^2}{\partial x'^2} + U(x', t) \right\} \mathcal{G}^{<}(x, x'; t) \end{split}$$

• Implemented via the Split Operator Method: Small $\Delta t \Rightarrow \mathcal{G}^{<}(t + \Delta t) \sim e^{-i\left\{\frac{\nabla^{2}}{2m} + U(x)\right\}\frac{\Delta t}{\hbar}}\mathcal{G}^{<}(t)e^{+i\left\{\frac{\nabla'^{2}}{2m} + U(x')\right\}\frac{\Delta t}{\hbar}}$ $e^{i(\hat{T}+\hat{U})\Delta t} \sim e^{i\frac{\hat{U}}{2}\Delta t}e^{i\hat{T}\Delta t}e^{i\frac{\hat{U}}{2}\Delta t} + O[\Delta t^{3}]$

Calculations in a box & FFT to switch representations

Mean-field evolution: implementation



· The mean-field is time-local

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$$\Sigma_{HF}(\mathbf{11}') = \delta(t_1 - t_{1'}) \Sigma_{HF}(x_1, x_{1'})$$

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Calculations in a box & FFT to switch representations

Mean-field TDGF vs. TDHF



- MF-TDGF and TDHF are numerically equivalent...
- but expressed in different terms!

 $i\frac{\partial}{\partial t}\mathcal{G}^{\leq}(x,x';t) = \left\{-\frac{1}{2m}\frac{\partial^2}{\partial x^2} + U(x)\right\}\mathcal{G}^{\leq}(x,x';t)$ $-\left\{-\frac{1}{2m}\frac{\partial^2}{\partial x'^2} + U(x')\right\}\mathcal{G}^{\leq}(x,x';t)$

Time Dependent Green's Functions

- 1 equation ... $N_x \times N_x$ matrix
- Testing ground
- Natural extension to correlated case via KB

Time Dependent Hartree-Fock

for
$$\alpha = 1, ..., N_{\alpha}$$

$$i \frac{\partial}{\partial t} \phi_{\alpha}(x, t) = \left\{ -\frac{1}{2m} \frac{\partial^2}{\partial x^2} + U(x) \right\} \phi_{\alpha}(x, t)$$

 \mathbf{end}

- N_α equations ... vectors of size N_x
- Limited to mean-field!
- Extension needs additional assumptions



- Initial state should be ground state of the Hamiltonian
 - Mean-field approx. \Rightarrow solve static Hartree-Fock equations
- Possible solution: use adiabatic theorem!

$$H(t) = f(t)H_0 + [1 - f(t)]H_1$$
$$f(t) = \begin{cases} 1, & t \to -\infty \\ 0, & t \to t_0 \end{cases}$$

- Advantage: a single code for everything!
- For practical applications:
 - $H_0 \& H_1$ with similar spectra to avoid crossing
 - $H_0 = \frac{1}{2}kx^2$
 - $H_1 = U_{\rm mf}$
 - Adiabatic transient: $f(t) = \frac{1}{1+e^{(t-\tau_0)/\tau}}, \qquad \tau \to \infty$



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Initial state and adiabatic switching



- Initial state should be ground state of the Hamiltonian
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Adiabatic switching & observables







Off-diagonal elements: origin





- $\mathcal{G}(x, x')$'s are matrices in x and x'
- Off-diagonal elements describe correlation of single-particle states $\mathcal{G}^<(x,x')=\sum_{\alpha=0}^{N_\alpha}\phi_\alpha(x)\phi^*_\alpha(x')$
- Diagonal elements yield physical properties $n(x) = \mathcal{G}^{<}(x, x' = x) = \sum_{\alpha=0}^{N_{\alpha}} n_{\alpha} |\phi_{\alpha}(x)|^{2} \qquad K = \sum_{k} \frac{k^{2}}{2m} \mathcal{G}^{<}(k, k' = k)$

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Off-diagonal elements: importance





Conceptual issues:

- Should far away sp states be connected in a nuclear reaction?
- Decoherence and dissipation will dominate late time evolution...
- Are $x \neq x'$ elements really necessary for the time-evolution?

Practical issues:

- Green's functions are $N^D_x \times N^D_x \times N^2_t$ matrices: $20^6 \sim 10^8$
- Eliminating off-diagonalities drastically reduces numerical cost

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Research program





A. Rios et al., in preparation.

- Used adiabatic theorem to solve mean-field $\sqrt{}$
- Full (N_x^2) , damped & cut $(N_a imes N_r)$ 1D mean-field evolution $\sqrt{}$
- Full 1D correlated evolution: Born approximation $\sim \sqrt{}$

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- Traditional calculations performed on $N_x \times N_x$ mesh
- Periodic boundary conditions
- Rotated coordinate frame: $x_a = \frac{x+x'}{2}, x_r = x' x$
- Control lengths and meshpoints $\Rightarrow (L_a, N_a) \times (L_r, N_r)$
- Reduce numerical effort by factors of 2-10





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Traditional vs. rotated evolutions





Two time Kadanoff-Baym equations



- Need of a strategy to deal with memory & two-times
- Time off-diagonal time elements are present
- Use symmetries $\mathcal{G}^{\lessgtr}(1,2) = -[\mathcal{G}^{\lessgtr}(2,1)]^*$ to minimize resources
- Self-consistency imposed at every time step





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Strategy to solve two-time equations





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Nuclear time-dependent correlations Questions & expectations



Some experience already gathered for uniform systems

Danielewicz, Ann. Phys. 152, 239 (1984)

H. S. Köhler, PRC 51 3232 (1995)

- Expected physical effects
 - Thermalization ($0 < n_{\alpha} < 1$)
 - Damping of collective modes
- · Correlations in the initial state
 - · Will a mean-field system evolve to a correlated ground state?
 - · Adiabatic switching on of correlations?
 - · Imaginary time evolution to get ground states?
- Testing ground calculations: 1D fermions on a HO trap
 - No mean-field, only confining potential
 - Test with mock gaussian NN force
 - Issues with cross section in 1D